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An Application of Laplace Transform, Inverse Laplace Transform and Pade's Approximant in the Periodic Solution of Duffing Equation of Motion

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ABSTRACT

The main objective of this research is to obtain the periodic solution by inserting the ideas of Laplace Transformation, Pade's Approximant and Inverse Laplace Transformation are applied. In this research article the Standard Duffing Equation of motion possessing symmetric oscillations is examined by using the Sumudu Transform Series Decomposition Method. The Duffing oscillator is a periodically forced oscillator with a nonlinear elasticity. For this type of oscillatory system there are frequencies at which the vibration suddenly jumps and these jumps depend upon whether the frequency is increasing or decreasing. Between these frequencies, multiple solutions exist for a given frequency of excitation, and the initial conditions determine which of these solutions represents the response of the system. The Sumudu Decomposition Transform Series Method (STSDM) is claimed as an efficient method for solving problems in Engineering and Science fields. It gives divergent series solution.

This article explores on studying the application capacity of STSDM in getting solution for symmetric oscillations of Duffing Equation of Motion which are simple and nonlinear.

Key words: Sumudu Transform (ST); Inverse Sumudu Transform (IST), Adomian Polynomials; Duffing Equation, Truly Nonlinear Oscillators, Laplace Transform (LT), Inverse Laplace Transform (ILT); Pade's Approximant (PA).

1. INTRODUCTION

A nonlinear second-order ordinary differential equation of the Duffing oscillator is [1-15].

$$\ddot{x} + \xi x + g(x) = G_0(t) \tag{1}$$

Over dot denote differentiation with respect to time, t. ξ is the damping factor. $G_0(t)$ is a time dependent forcing function. The cubic polynomial restoring force function is

$$g(x) = ax + bx^2 + cx^3 + d$$
 (2)

Eq. (2) represents:

Asymmetric oscillations for $b \neq 0$ and/or $d \neq 0$. Symmetric oscillations for b = d = 0.

c > 0, is for the hardening type of the system.

c < 0, is for the softening type of the system.

STSDM is considered to be an efficient tool for cracking a large number of physical problems in Science and Engineering. Sumudu Transform reduces the complexity in the integration of highly integral functions which are nonlinear. The rate of convergence increases the solution's series expansion. Nonlinear terms in the DE are resolved by the Adomian Polynomials.

This paper studies the application capacity of STSDM in getting solution for symmetric oscillations of Duffing Equation of Motion which are simple and nonlinear.

2. ANALYSIS

A Duffing Equation of Motion which is simple and nonlinear is obtained by

$$\dot{x} + x + 0.3x^3 = 0$$
 (3)
 $x = 0$ and $\dot{x} = 1$ at $t = 0$ (4)

the restoring force function is $g(x) = x + 0.3x^3$.

If (3) and (4) are applied by ST[18-23], one can see

$$S\{x(t)\} = v - v^2 S[x(t) + 0.3x^3(t)]$$
(5)

If (5) is applied by IST, the result becomes

$$x(t) = t - S^{-1} \Big[v^2 S \Big(x(t) + 0.3 x^3(t) \Big) \Big]$$
(6)

Assuming the series solution $x(t) = \sum_{m=0}^{\infty} x_m(t)$ for

Eq.(3) and expressing (6) as

$$\sum_{m=0}^{\infty} x_m(t) = t - S^{-1} \left[v^2 S \left(\sum_{m=0}^{\infty} A_m \right) \right]$$
(7)

The Adomian Polynomial Functions in (7) are

$$A_{m} = \frac{1}{m!} \frac{d^{m}}{d\theta^{m}} \left[\sum_{j=0}^{\infty} \theta^{j} x_{j} + 0.3 \left(\sum_{j=0}^{\infty} \theta^{j} x_{j} \right)^{3} \right] \Big|_{\theta=0}$$
$$= x_{m} + 0.3 \sum_{j=0}^{\infty} x_{j} x_{m-j}$$
(8)

Comparing like terms on both sides and using (8) in (7) one can have

$$x_0(t) = t \tag{9}$$

$$x_{m+1}(t) = -S^{-1} \Big[v^2 S \Big(A_m \Big) \Big] \quad m \ge 0$$
 (10)

$$x_1(t) = -\frac{t^3}{6}$$
(11)

$$x_2(t) = -0.8 \frac{t^5}{120} \tag{12}$$

$$x_3(t) = 18.8 \frac{t^7}{5040} \tag{13}$$

the series solution of (3) with boundary constraints (4) are got as

$$x(t) = \sum_{m=0}^{\infty} x_m(t)$$
(14)
(15)

$$x(t) = t - \frac{1}{6} - 0.8 \frac{1}{120} + 18.8 \frac{1}{5040} + \dots$$
 (15)
The series solution (15) is unable to exhibit

The series solution (15) is unable to exhibit the periodicity.

Applying LT, $\begin{bmatrix} 4\\4 \end{bmatrix}$ PA, and the ILT, the solution of the problem is got as

 $x(t) = 0.928746 \ SIN (1.10982 \ t) - 0.0103828 \ SIN (2.96113 \ t)$ (16) The phase diagram for (1) with the boundary constraints (2) arises as

$$(\dot{x})^2 = 1 - x^2 - 0.15x^4$$

$$= (1 + 0.13246x^2)(1 - 1.13246x^2)$$
(17)

For +ve and -ve amplitudes the behavior oscillations are one and the same. The singular point of DE(1) in the phase diagram (that is \dot{x} Vs x curve) got from the zeros of g(x) is (0,0). The primitive of g(x) w.r.t is x is $g'(x) = 1 + 0.3x^2$ and g(0)=1>0, which implies that the singular point

(0,0) is turned into centre. If $g'(x^*)$ is negative

then x^* is a saddle point.

Eq.(17) gives the phase diagram (\dot{x} Vs x) for the DE (1) with boundary constraints (2). \dot{x} Vs x plot arises out of Eq.(17) and this graph depicts the boundary which is closed. Consequently the existence of the periodic solution is being obtained. Eq. (17) depicts the equal magnitude of +ve & -ve amplitudes (that is x=+0.9397 or -.09397, \dot{x} =0), when x=0, \dot{x} = +1 or \dot{x} = -1. From the STSDM solution (16) one can have \dot{x} values as 1.047,-1.016 at x=0. The STSDM gives the solution exactly in the neighborhood of the domain in which boundary constraints are described.



Figure 1 :Phase diagrams of Duffing equation arised out of the solution of STSDM versus Exact Solution

The solution obtained by the STSDM method is divergent in order to make the solution convergent the concept of Laplace transformation and Pade's approximant and Inverse laplace transformation is applied like the procedure followed for the modified differential transform method.

Praveen & Rao [10], solved the duffing equation of motion having the symmetric oscillations. The solution obtained by the above mentioned procedure gives the solution exactly near the region where the initial conditions are specified. They used the Harmonic balance method to examine the performance of the nonlinear oscillations. The solution obtained by using the Harmonic balance method coincides exact solution.

The following diagram shows the comparison of the solution obtained by Harmonic balance method with the exact solution.



Figure 2 :Phase diagrams of Duffing equation arises out of the solution of Harmonic balance method Vs Exact Solution

The discrepancy in solution (16) is mainly due to the first harmonic frequency and is not three times of third harmonic frequency, which results in different magnitudes. In their study they also suggested some modification to the procedure and obtained the solution and again compared the solution with the exact solution which is able to capture the actual trend in the phase diagram.

The following figure 3 shows the comparison of the solution obtained after modification to the solution obtained by MDTM to the exact solution.



Figure 3: Phase diagrams of Duffing equation arised out of the solution of modified STSDM solution versus Exact Solution

3. CONCLUSION

The behaviour of oscillations of the Duffing equation of motion (1) with boundary constraints (2) can easily be understood through generation of phase diagrams. Since the controlling equation is a second order NLODE, one has to verify the approximate solution and its first derivative with the numerical solution. Hence in this research article, the first order differential equation is arrived after integrating the Duffing equation of motion and applying the boundary constarints, which can depicts the exact phase diagram. The periodic solution of the Duffing equation of motion obtained by STSDM is not capturing equal magnitudes of first derivative of x at x=0, t=0 and at the period of oscillations. The STSDM provides

the solution accurately near to the region where the initial conditions are specified. It follows the actual trend in the phase diagram when the solution (16) is modified.

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