



Coupled Dynamic Thermoplasticity Problem for Transversely Isotropic Parallelepiped

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ABSTRACT

The article formulates a coupled thermoplastic dynamic boundary value problem based on the deformation theory of transversally isotropic bodies. The coupled boundary value problem consists of the motion equation, the thermoplasticity constitutive relations for transversely isotropic bodies, the Cauchy relation and the heat conduction equation, with the corresponding initial and boundary conditions. A coupled dynamic boundary value problem for a thermoplastic transversely isotropic parallelepiped is considered. For the numerical solution of the coupled boundary value problem, an explicit finite-difference schemes are constructed. The finite-difference equations are written in the form of recurrent relations allowing to find the components of displacement and temperature, taking into account the initial and boundary conditions at each layer in time. On the basis of the obtained numerical results, the stress-strain state of a transversely isotropic thermoplastic parallelepiped is investigated and the propagation of the plastic zone is shown.

Key words: Coupled problem, thermoplasticity, temperature, displacements, difference scheme, plastic zone, deformation, stress.

1. INTRODUCTION

Investigation the joint influence of the thermomechanical forces on the strength and durability of structures is an actual problem of engineering. Usually this type of problems in the framework of the thermoelasticity or thermoplasticity can be described by the system of motion and heat conduction equations with a corresponding initial and thermo-mechanical boundary conditions.

The coupled thermoelastic boundary value problems was investigated firstly by Biot [1]. In the works of Lord & Shulman[2] has been introduced a generalized coupled theory with a wave-type heat equation The following works [3-10] are devoted to the development of the theoretical foundations of coupled thermomechanical deformations of solids subject to both large and inelastic deformations. The coupled

thermoplasticity problems are considered in [7,9-11,15,16]. The coupled and uncoupled thermomechanical boundary value problems are numerically solved in following works [12-14, 17-18,20].

The paper [18] considers the statement and numerical solution of partially related problems of thermoelasticity. In work [19] the analysis of heat transfer of water heating convectors is given.

In Sections~2 using the deformation thermoplasticity theory for transversely isotropic materials the dynamic coupled boundary value problem are formulated. Note, the boundary value problems, are reduced to the system consisting of three nonlinear motion and one heat partial differential equations.

In Section~3, using the finite difference method, the discrete analogy of the coupled dynamic boundary value problems based on deformation theory of thermoplasticity are constructed. The finite-difference equations resolved in relative to the desired quantities are reduced to recurrent formulas and taking into account the initial and boundary conditions a numerical results are determined.

In section~4, a clamped on all surfaces parallelepiped with a given initial sinusoidal thermal load is considered. The described process using the deformation theory of thermoplasticity is formulated as a coupled boundary value problem and numerically solved. The temperature distribution and the appearing of the plasticity zones inside the parallelepiped under the given initial and boundary conditions are investigated.

2. FORMULATION OF THE COUPLED DYNAMIC THERMOPLASTICITY BOUNDARY VALUE PROBLEM

Consider a coupled dynamic boundary value problem of thermoplasticity for a transversely isotropic parallelepiped consisting of the motion equation

$$\sum_{j=1}^3 \frac{\partial \sigma_{ij}}{\partial x_j} + X_i = \rho \frac{\partial^2 u_i}{\partial t^2}, \quad i = 1, 2, 3 \quad (1)$$

constitutive relation of the deformation theory of thermoplasticity for transversely isotropic bodies [17,20]

$$\sigma_{ij} = \begin{cases} C_{ijk}\varepsilon_{kl} - \beta_{ij}(T-T_0)\delta_{ij}, & \text{if } p < p^*, q < q^* \\ C_{ijk}\varepsilon_{kl} - \beta_{ij}(T-T_0)\delta_{ij} - 2(\lambda_2 - \lambda_2')(1 - \frac{p^*}{p})p_{ij} - \\ - 2(\lambda_5 - \lambda_5')(1 - \frac{q^*}{q})q_{ij}, & \text{if } p \geq p^*, q \geq q^* \end{cases} \quad (2)$$

Cauchy relations

$$\varepsilon_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right), \quad (3)$$

heat conduction equations for anisotropic bodies

$$\lambda_{ij}T_{,ij} - c_\varepsilon \dot{T} - T \cdot \beta_{ij} \cdot \dot{\varepsilon}_{ij} = 0 \quad (4)$$

with initial

$$u_i|_{t=t_0} = \varphi_i, \quad \dot{u}_i|_{t=t_0} = \psi_i, \quad T|_{t=t_0} = T_0 \quad (5)$$

and boundary conditions

$$u_i|_{\Sigma_1} = u_i^0, \quad T|_{\Sigma} = \bar{T}_0, \quad \sum_{j=1}^3 \sigma_{ij}n_j|_{\Sigma_2} = S_i \quad (6)$$

where, σ_{ij} - stress tensor, ε_{ij} - strain tensor, u_i - displacement components, X_i - volume force, C_{ijkl} - tensor of the fourth rank determining the mechanical properties of a material, T -temperature, T_0 -initial temperature, p, q - strain tensor intensity, p^*, q^* - elastic limits in the longitudinal and transverse directions of a transversely isotropic body, β_{ij} - corresponds to thermal expansion coefficients, ρ -density of the body, δ_{ij} - delta Kronecker symbol, n_j - external normal to the surface Σ_2 , S_1, S_2, S_3 - components of the external load vector, λ_{ij} - is the heat flow coefficients, c_ε -denotes heat coefficient at a constant deformation, $\Sigma = \Sigma_1 + \Sigma_2$ - the surface of the body under consideration, the upper point denotes the time derivative.

Equations (1)-(6) represent a dynamic coupled boundary value problem of thermoplasticity for transversely isotropic bodies. The boundary value problem (1)-(6), after some transformations, can be written in the following form, with respect to displacements and temperature

$$C_{1111} \frac{\partial^2 u}{\partial x^2} + C_{1212} \frac{\partial^2 u}{\partial y^2} + C_{1313} \frac{\partial^2 u}{\partial z^2} + (C_{1122} + C_{1212}) \frac{\partial^2 v}{\partial x \partial y} + (C_{1133} + C_{1313}) \frac{\partial^2 w}{\partial x \partial z} - \beta_{11} \frac{\partial T}{\partial x} - F_1 = \rho \quad (7)$$

$$C_{1212} \frac{\partial^2 v}{\partial x^2} + C_{2222} \frac{\partial^2 v}{\partial y^2} + C_{2323} \frac{\partial^2 v}{\partial z^2} + (C_{2211} + C_{1212}) \frac{\partial^2 u}{\partial x \partial y} + (C_{2233} + C_{2323}) \frac{\partial^2 w}{\partial y \partial z} - \beta_{22} \frac{\partial T}{\partial y} - F_2 = \rho \frac{\partial^2 v}{\partial t^2} \quad (8)$$

$$C_{1313} \frac{\partial^2 w}{\partial x^2} + C_{2323} \frac{\partial^2 w}{\partial y^2} + C_{3333} \frac{\partial^2 w}{\partial z^2} + (C_{3311} + C_{1313}) \frac{\partial^2 u}{\partial x \partial z} + (C_{3322} + C_{2323}) \frac{\partial^2 v}{\partial y \partial z} - \beta_{33} \frac{\partial T}{\partial z} - F_3 = \rho \frac{\partial^2 w}{\partial t^2} \quad (9)$$

heat flow equation

$$\lambda_{11} \frac{\partial^2 T}{\partial x^2} + \lambda_{22} \frac{\partial^2 T}{\partial y^2} + \lambda_{33} \frac{\partial^2 T}{\partial z^2} - c_\varepsilon \frac{\partial T}{\partial t} - T \left(\beta_{11} \frac{\partial^2 u}{\partial x \partial t} + \beta_{22} \frac{\partial^2 v}{\partial y \partial t} + \beta_{33} \frac{\partial^2 w}{\partial z \partial t} \right) = 0 \quad (10)$$

with corresponding initial

$$u(x, y, z, t)|_{t=0} = \varphi_1, \quad v(x, y, z, t)|_{t=0} = \varphi_2, \quad w(x, y, z, t)|_{t=0} = \varphi_3,$$

$$\frac{\partial u}{\partial t}|_{t=0} = \psi_1, \quad \frac{\partial v}{\partial t}|_{t=0} = \psi_2, \quad \frac{\partial w}{\partial t}|_{t=0} = \psi_3,$$

$$T(x, y, z, t)|_{t=0} = T_0$$

and boundary conditions with respect to displacements

$$u(x, y, z, t)|_{\Sigma_1} = u^0, \quad v(x, y, z, t)|_{\Sigma_1} = v^0, \quad w(x, y, z, t)|_{\Sigma_1} = w^0$$

or loads

$$(\sigma_{11}n_1 + \sigma_{12}n_2 + \sigma_{13}n_3)|_{\Sigma_2} = S_1, \quad (\sigma_{21}n_1 + \sigma_{22}n_2 + \sigma_{23}n_3)|_{\Sigma_2} = S_2,$$

$$(\sigma_{31}n_1 + \sigma_{32}n_2 + \sigma_{33}n_3)|_{\Sigma_2} = S_3$$

and temperatures

$$T(x, y, z, t)|_{\Sigma} = \bar{T}_0$$

where

$$F_1 = 2(\lambda_2 - \lambda_2')(1 - \frac{p^*}{p}) \left[\frac{\partial(\varepsilon_{11} - \varepsilon_{22})}{2\partial x_1} + \frac{\partial \varepsilon_{12}}{\partial x_2} \right] + 2(\lambda_5 - \lambda_5')(1 - \frac{p^*}{p}) \frac{\partial \varepsilon_{13}}{\partial x_3},$$

$$F_2 = 2(\lambda_2 - \lambda_2')(1 - \frac{p^*}{p}) \left[\frac{\partial \varepsilon_{21}}{\partial x_2} + \frac{\partial(\varepsilon_{22} - \varepsilon_{11})}{2\partial x_2} \right] + 2(\lambda_5 - \lambda_5')(1 - \frac{p^*}{p}) \frac{\partial \varepsilon_{23}}{\partial x_3},$$

$$F_3 = 2(\lambda_5 - \lambda_5')(1 - \frac{q^*}{q}) \left[\frac{\partial \varepsilon_{31}}{\partial x_1} + \frac{\partial \varepsilon_{32}}{\partial x_2} \right].$$

if $p \geq p^*, q \geq q^*$.

3. THE FINITE DIFFERENCE EQUATIONS FOR COUPLED DYNAMIC BOUNDARY VALUE PROBLEMS OF THERMO PLASTISITY

The coupled thermoplasticity boundary value problem based on deformation theory of thermoplasticity for transversely isotropic materials presents by the eqs. (7-10) considered in the domain $\Omega = \{t \geq 0, 0 \leq x \leq l_1, 0 \leq y \leq l_2, 0 \leq z \leq l_3\}$.

In order to construct a grid equations consider three sets of parallel lines $x_i = ih_1$ ($i=0, N_1$), $y_j = jh_2$ ($j=0, N_2$),

$z_k = kh_3$ ($k=0, N_3$) where $h_1 = l_1 / N_1, h_2 = l_2 / N_2,$

$h_3 = l_3 / N_3$, in the domain Ω and taking

$t_m = m\tau$ ($m=0, 1, 2, \dots$), where τ - a step on the axis t .

Then replacing the derivatives in eqs.(7-10) by difference quotients, we obtain

$$\begin{aligned}
 & C_{1111} \frac{u_{i+1,j,k}^{n+1} - 2u_{i,j,k}^{n+1} + u_{i-1,j,k}^{n+1}}{h_1^2} + C_{1212} \frac{u_{i,j+k}^{n+1} - 2u_{i,j,k}^{n+1} + u_{i,j-1,k}^{n+1}}{h_2^2} + \\
 & + C_{1313} \frac{u_{i,j,k+1}^{n+1} - 2u_{i,j,k}^{n+1} + u_{i,j,k-1}^{n+1}}{h_3^2} + (C_{1122} + C_{1212}) \frac{v_{i+1,j+1,k}^{n+1} - v_{i-1,j+1,k}^{n+1} - v_{i+1,j-1,k}^{n+1} + v_{i-1,j-1,k}^{n+1}}{4h_1h_2} + \\
 & + (C_{1133} + C_{1313}) \frac{w_{i+1,j,k+1}^{n+1} - w_{i-1,j,k+1}^{n+1} - w_{i+1,j,k-1}^{n+1} + w_{i-1,j,k-1}^{n+1}}{4h_1h_3} - \\
 & - \beta_{11} \frac{T_{i+1,j,k}^{n+1} - T_{i-1,j,k}^{n+1}}{2h_1} - F_1 = \rho \frac{u_{i,j,k}^{n+1} - 2u_{i,j,k}^{n+1} + u_{i,j,k}^{n-1}}{\tau^2}
 \end{aligned} \tag{11}$$

$$\begin{aligned}
 & C_{1212} \frac{v_{i+1,j,k}^{n+1} - 2v_{i,j,k}^{n+1} + v_{i-1,j,k}^{n+1}}{h_1^2} + C_{2222} \frac{v_{i,j+k}^{n+1} - 2v_{i,j,k}^{n+1} + v_{i,j-1,k}^{n+1}}{h_2^2} + \\
 & + C_{2323} \frac{v_{i,j,k+1}^{n+1} - 2v_{i,j,k}^{n+1} + v_{i,j,k-1}^{n+1}}{h_3^2} + (C_{2111} + C_{1212}) \frac{u_{i+1,j+1,k}^{n+1} - u_{i-1,j+1,k}^{n+1} - u_{i+1,j-1,k}^{n+1} + u_{i-1,j-1,k}^{n+1}}{4h_1h_2} + \\
 & + (C_{2233} + C_{2323}) \frac{w_{i,j+1,k+1}^{n+1} - w_{i,j-1,k+1}^{n+1} - w_{i,j+1,k-1}^{n+1} + w_{i,j-1,k-1}^{n+1}}{4h_2h_3} - \\
 & - \beta_{22} \frac{T_{i,j+1,k}^{n+1} - T_{i,j-1,k}^{n+1}}{2h_2} - F_2 = \rho \frac{v_{i,j,k}^{n+1} - 2v_{i,j,k}^{n+1} + v_{i,j,k}^{n-1}}{\tau^2}
 \end{aligned} \tag{12}$$

$$\begin{aligned}
 & C_{1313} \frac{w_{i+1,j,k}^{n+1} - 2w_{i,j,k}^{n+1} + w_{i-1,j,k}^{n+1}}{h_1^2} + C_{2323} \frac{w_{i,j+k}^{n+1} - 2w_{i,j,k}^{n+1} + w_{i,j-1,k}^{n+1}}{h_2^2} + \\
 & + C_{3333} \frac{w_{i,j,k+1}^{n+1} - 2w_{i,j,k}^{n+1} + w_{i,j,k-1}^{n+1}}{h_3^2} + (C_{3311} + C_{1313}) \frac{u_{i+1,j,k+1}^{n+1} - u_{i-1,j,k+1}^{n+1} - u_{i+1,j,k-1}^{n+1} + u_{i-1,j,k-1}^{n+1}}{4h_1h_3} + \\
 & + (C_{3322} + C_{2323}) \frac{v_{i,j+1,k+1}^{n+1} - v_{i,j-1,k+1}^{n+1} - v_{i,j+1,k-1}^{n+1} + v_{i,j-1,k-1}^{n+1}}{4h_2h_3} - \\
 & - \beta_{33} \frac{T_{i,j,k+1}^{n+1} - T_{i,j,k-1}^{n+1}}{2h_3} - F_3 = \rho \frac{w_{i,j,k}^{n+1} - 2w_{i,j,k}^{n+1} + w_{i,j,k}^{n-1}}{\tau^2}
 \end{aligned} \tag{13}$$

$$\begin{aligned}
 & \lambda_{11} \frac{T_{i+1,j,k}^{n+1} - 2T_{i,j,k}^{n+1} + T_{i-1,j,k}^{n+1}}{h_1^2} + \lambda_{22} \frac{T_{i,j+k}^{n+1} - 2T_{i,j,k}^{n+1} + T_{i,j-1,k}^{n+1}}{h_2^2} + \lambda_{33} \frac{T_{i,j,k+1}^{n+1} - 2T_{i,j,k}^{n+1} + T_{i,j,k-1}^{n+1}}{h_3^2} - \\
 & - c_e \frac{T_{i,j,k}^{n+1} - T_{i,j,k}^{n-1}}{\tau} - T_{ij} \left(\beta_{11} \frac{u_{i+1,j,k+1}^{n+1} - u_{i-1,j,k+1}^{n+1} - u_{i+1,j,k-1}^{n+1} + u_{i-1,j,k-1}^{n+1}}{4h_1\tau} + \right. \\
 & \left. + \beta_{22} \frac{v_{i,j+1,k+1}^{n+1} - v_{i,j-1,k+1}^{n+1} - v_{i,j+1,k-1}^{n+1} + v_{i,j-1,k-1}^{n+1}}{4h_2\tau} + \beta_{33} \frac{w_{i,j,k+1}^{n+1} - w_{i,j,k-1}^{n+1} - w_{i,j,k+1}^{n-1} + w_{i,j,k-1}^{n-1}}{4h_3\tau} \right) = 0
 \end{aligned} \tag{14}$$

Solving the difference equations (11-14), with respect to $u_{i,j,k}^{n+1}$, $v_{i,j,k}^{n+1}$, $w_{i,j,k}^{n+1}$ and $T_{i,j,k}^{n+1}$, accordingly, we obtain

$$\begin{aligned}
 & u_{i,j,k}^{n+1} = \frac{\tau^2}{\rho} \left(C_{1111} \frac{u_{i+1,j,k}^n - 2u_{i,j,k}^n + u_{i-1,j,k}^n}{h_1^2} + C_{1212} \frac{u_{i,j+k}^n - 2u_{i,j,k}^n + u_{i,j-1,k}^n}{h_2^2} + \right. \\
 & + C_{1313} \frac{u_{i,j,k+1}^n - 2u_{i,j,k}^n + u_{i,j,k-1}^n}{h_3^2} + (C_{1122} + C_{1212}) \frac{v_{i+1,j+1,k}^n - v_{i-1,j+1,k}^n - v_{i+1,j-1,k}^n + v_{i-1,j-1,k}^n}{4h_1h_2} + \\
 & + (C_{1133} + C_{1313}) \frac{w_{i+1,j,k+1}^n - w_{i-1,j,k+1}^n - w_{i+1,j,k-1}^n + w_{i-1,j,k-1}^n}{4h_1h_3} - \\
 & - \beta_{11} \frac{T_{i+1,j,k}^n - T_{i-1,j,k}^n}{2h_1} - F_1 \left. \right) + 2u_{i,j,k}^n - u_{i,j,k}^{n-1}
 \end{aligned} \tag{15}$$

$$\begin{aligned}
 & v_{i,j,k}^{n+1} = \frac{\tau^2}{\rho} \left(C_{1212} \frac{v_{i+1,j,k}^n - 2v_{i,j,k}^n + v_{i-1,j,k}^n}{h_1^2} + C_{2222} \frac{v_{i,j+k}^n - 2v_{i,j,k}^n + v_{i,j-1,k}^n}{h_2^2} + \right. \\
 & + C_{2323} \frac{v_{i,j,k+1}^n - 2v_{i,j,k}^n + v_{i,j,k-1}^n}{h_3^2} + (C_{2111} + C_{1212}) \frac{u_{i+1,j+1,k}^n - u_{i-1,j+1,k}^n - u_{i+1,j-1,k}^n + u_{i-1,j-1,k}^n}{4h_1h_2} + \\
 & + (C_{2233} + C_{2323}) \frac{w_{i,j+1,k+1}^n - w_{i,j-1,k+1}^n - w_{i,j+1,k-1}^n + w_{i,j-1,k-1}^n}{4h_2h_3} - \\
 & - \beta_{22} \frac{T_{i,j+1,k}^n - T_{i,j-1,k}^n}{2h_2} - F_2 \left. \right) + 2v_{i,j,k}^n - v_{i,j,k}^{n-1}
 \end{aligned} \tag{16}$$

$$\begin{aligned}
 & w_{i,j,k}^{n+1} = \frac{\tau^2}{\rho} \left(C_{1313} \frac{w_{i+1,j,k}^n - 2w_{i,j,k}^n + w_{i-1,j,k}^n}{h_1^2} + C_{2323} \frac{w_{i,j+k}^n - 2w_{i,j,k}^n + w_{i,j-1,k}^n}{h_2^2} + \right. \\
 & + C_{3333} \frac{w_{i,j,k+1}^n - 2w_{i,j,k}^n + w_{i,j,k-1}^n}{h_3^2} + (C_{3311} + C_{1313}) \frac{u_{i+1,j,k+1}^n - u_{i-1,j,k+1}^n - u_{i+1,j,k-1}^n + u_{i-1,j,k-1}^n}{4h_1h_3} + \\
 & + (C_{3322} + C_{2323}) \frac{v_{i,j+1,k+1}^n - v_{i,j-1,k+1}^n - v_{i,j+1,k-1}^n + v_{i,j-1,k-1}^n}{4h_2h_3} - \\
 & - \beta_{33} \frac{T_{i,j,k+1}^n - T_{i,j,k-1}^n}{2h_3} - F_3 \left. \right) + 2w_{i,j,k}^n - w_{i,j,k}^{n-1}
 \end{aligned} \tag{17}$$

$$\begin{aligned}
 & T_{i,j,k}^{n+1} = \frac{\tau}{c_e} \left(\lambda_{11} \frac{T_{i+1,j,k}^n - 2T_{i,j,k}^n + T_{i-1,j,k}^n}{h_1^2} + \lambda_{22} \frac{T_{i,j+k}^n - 2T_{i,j,k}^n + T_{i,j-1,k}^n}{h_2^2} + \lambda_{33} \frac{T_{i,j,k+1}^n - 2T_{i,j,k}^n + T_{i,j,k-1}^n}{h_3^2} + \right. \\
 & + \lambda_{33} \frac{T_{i,j,k+1}^n - 2T_{i,j,k}^n + T_{i,j,k-1}^n}{h_3^2} - T_{ij} \left(\beta_{11} \frac{u_{i+1,j,k+1}^{n+1} - u_{i-1,j,k+1}^{n+1} - u_{i+1,j,k-1}^{n+1} + u_{i-1,j,k-1}^{n+1}}{4h_1\tau} + \right. \\
 & \left. + \beta_{22} \frac{v_{i,j+1,k+1}^{n+1} - v_{i,j-1,k+1}^{n+1} - v_{i,j+1,k-1}^{n+1} + v_{i,j-1,k-1}^{n+1}}{4h_2\tau} + \beta_{33} \frac{w_{i,j,k+1}^{n+1} - w_{i,j,k-1}^{n+1} - w_{i,j,k+1}^{n-1} + w_{i,j,k-1}^{n-1}}{4h_3\tau} \right) \left. \right) + T_{i,j,k}^n
 \end{aligned} \tag{18}$$

The resulting recurrent formulas (15-18), taking into account the initial and boundary conditions, make it possible to find the numerical values of the sought functions at each layer in time. First, we solve the thermoelastic problem; in this case, the values of F_1, F_2, F_3 in formulas (15)-(17) are trivial. On each layer, plastic zones are determined in time, where the values of the sought functions are re-calculated.

4. NUMERICAL TEST

Let's consider a clamped all sides parallelepiped with an internal initial sinusoidal temperature at the time $t = 0$. The described process using the deformation thermoplasticity theory has been modelled by the eqs. (7-10). The initial conditions take the following forms:

$$\begin{aligned}
 & u(x, y, z, t) \Big|_{t=0} = 0, \quad v(x, y, z, t) \Big|_{t=0} = 0, \quad w(x, y, z, t) \Big|_{t=0} = 0 \\
 & \frac{\partial u}{\partial t} \Big|_{t=0} = 0, \quad \frac{\partial v}{\partial t} \Big|_{t=0} = 0, \quad \frac{\partial w}{\partial t} \Big|_{t=0} = 0,
 \end{aligned} \tag{19}$$

$$T(x, y, z, t) \Big|_{t=0} = T_0 \sin(\pi x_i) \sin(\pi y_j) \sin(\pi z_k)$$

and boundary conditions

$$\begin{aligned}
 & u(x, y, z, t) \Big|_{\Sigma} = 0, \quad v(x, y, z, t) \Big|_{\Sigma} = 0, \quad w(x, y, z, t) \Big|_{\Sigma} = 0, \\
 & T(x, y, z, t) \Big|_{\Sigma} = 0.
 \end{aligned} \tag{20}$$

Elastic moduli, hardening moduli, elastic limits, heat capacity at constant deformation, thermal expansion tensor, thermal conductivity tensor, density had the following dimensionless values:

$$C_{1122}=0.21, C_{1133}=0.19, C_{2233}=0.19, C_{3333}=5.356,$$

$$C_{2323}=2.39, C_{1313}=2.39,$$

$$C_{1212}=2.735, C_{1111}=5.68, C_{2222}=5.68, b_{11}=0.25, b_{22}=b_{11}, b_{33}=0.37,$$

$$\lambda_{11}=0.03, \lambda_{22}=0.03, \lambda_{33}=0.01, \lambda_2 = 2.5, \lambda_5 = 2.24, \rho=0.81,$$

$$c_e=3.5, \lambda_2 = 2.5, \lambda_5 = 2.24, \tau=0.01,$$

$$p^* = 0.08, q^* = 0.04, T_0 = 20, N_1 = N_2 = N_3 = 10.$$

Below we give the zones of plasticity in different layers in time and in sections along the coordinate axes of the considered transversely isotropic parallelepiped.

Table 1: Strain tensor intensity values p , at $z=0.7, t=0.07$

| $\begin{matrix} x \\ y \end{matrix}$ | 0 | 0.1 | 0.2 | 0.3 | 0.4 | 0.5 |
|--------------------------------------|---|---------|---------|---------|---------|---------|
| 0 | 0 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 |
| 0.1 | 0 | 0.10273 | 0.11462 | 0.12015 | 0.11984 | 0.11940 |
| 0.2 | 0 | 0.11439 | 0.08931 | 0.07044 | 0.04453 | 0.02822 |
| 0.3 | 0 | 0.12001 | 0.07069 | 0.05097 | 0.02713 | 0.00359 |
| 0.4 | 0 | 0.11963 | 0.04476 | 0.02715 | 0.01420 | 0.00012 |
| 0.5 | 0 | 0.11914 | 0.02836 | 0.00372 | 0.00027 | 0.00008 |
| 0.6 | 0 | 0.11960 | 0.04472 | 0.02715 | 0.01420 | 0.00012 |
| 0.7 | 0 | 0.11995 | 0.07066 | 0.05097 | 0.02713 | 0.00359 |
| 0.8 | 0 | 0.11432 | 0.08935 | 0.07055 | 0.04465 | 0.02822 |
| 0.9 | 0 | 0.10289 | 0.11478 | 0.12024 | 0.11987 | 0.11940 |
| 1 | 0 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 |

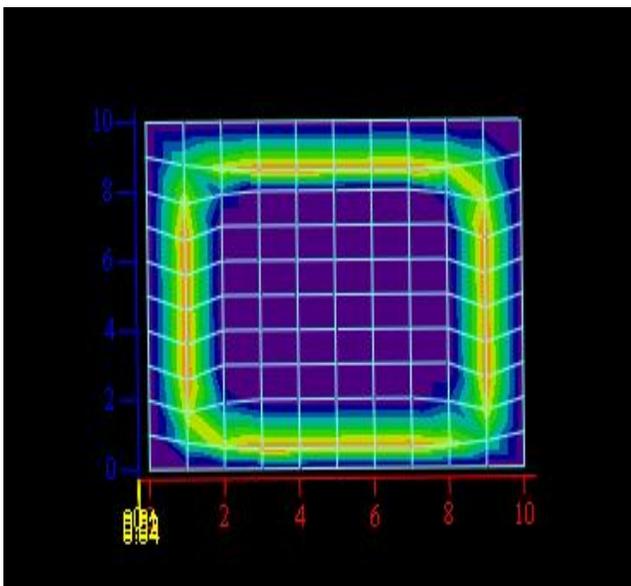


Figure 1: Plasticity zone by intensity tensor of deformations p in the XOY plane at $z=0.7, t=0.07 (p \geq p^*), p^*=0.08$

Table 2: Strain tensor intensity values q , at $y=0.9, t=0.07$

| $\begin{matrix} x \\ z \end{matrix}$ | 0 | 0.1 | 0.2 | 0.3 | 0.4 | 0.5 |
|--------------------------------------|---|---------|---------|---------|---------|---------|
| 0 | 0 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 |
| 0.1 | 0 | 0.05295 | 0.07524 | 0.09512 | 0.10844 | 0.11311 |
| 0.2 | 0 | 0.04995 | 0.07114 | 0.09020 | 0.10297 | 0.10744 |
| 0.3 | 0 | 0.03817 | 0.05432 | 0.06887 | 0.07861 | 0.08203 |
| 0.4 | 0 | 0.02026 | 0.02880 | 0.03651 | 0.04166 | 0.04346 |
| 0.5 | 0 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 |
| 0.6 | 0 | 0.02026 | 0.02880 | 0.03651 | 0.04166 | 0.04346 |
| 0.7 | 0 | 0.03817 | 0.05432 | 0.06887 | 0.07861 | 0.08203 |
| 0.8 | 0 | 0.04995 | 0.07114 | 0.09020 | 0.10297 | 0.10744 |
| 0.9 | 0 | 0.05295 | 0.07524 | 0.09512 | 0.10844 | 0.11311 |
| 1 | 0 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 |

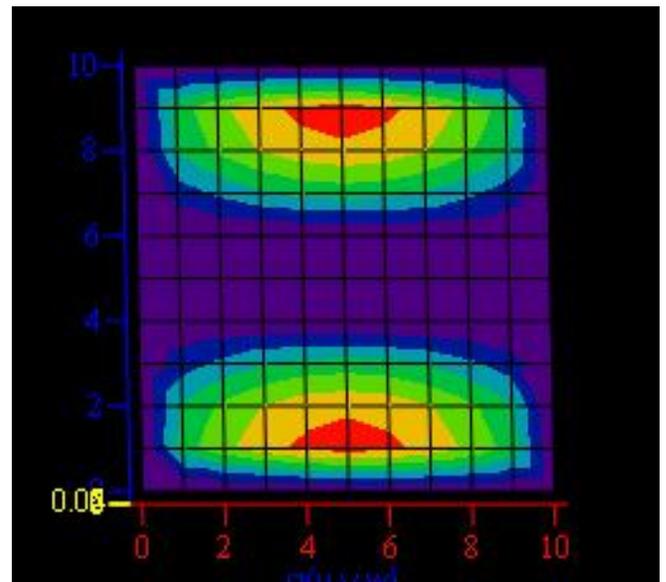


Figure 2: Plasticity zone by intensity of tensor deformations q in the XOY plane at $z=0.8, t=0.07 (q \geq q^*), q^*=0.04$

Next, we present the numerical results and 3D graphs of the functions $u(x, y, z, t), v(x, y, z, t), w(x, y, z, t)$ and $T(x, y, z, t)$ in the area under consideration.

Table 3: Displacement values $u(x,y,z,t)$ at $z=0.3$ and $t=0.09$

| $\begin{matrix} x \\ y \end{matrix}$ | 0 | 0.1 | 0.2 | 0.3 | 0.4 | 0.5 |
|--------------------------------------|---|----------|----------|----------|----------|---------|
| 0 | 0 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 |
| 0.1 | 0 | -0.01143 | -0.01372 | -0.01076 | -0.00574 | 0.00000 |
| 0.2 | 0 | -0.01898 | -0.02335 | -0.01833 | -0.00976 | 0.00000 |
| 0.3 | 0 | -0.02558 | -0.03158 | -0.02482 | -0.01322 | 0.00000 |
| 0.4 | 0 | -0.03002 | -0.03706 | -0.02913 | -0.01551 | 0.00000 |
| 0.5 | 0 | -0.03158 | -0.03896 | -0.03062 | -0.01631 | 0.00000 |
| 0.6 | 0 | -0.03005 | -0.03706 | -0.02913 | -0.01551 | 0.00000 |
| 0.7 | 0 | -0.02562 | -0.03159 | -0.02482 | -0.01322 | 0.00000 |
| 0.8 | 0 | -0.01903 | -0.02335 | -0.01833 | -0.00976 | 0.00000 |
| 0.9 | 0 | -0.01147 | -0.01372 | -0.01072 | -0.00569 | 0.00000 |
| 1 | 0 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 |

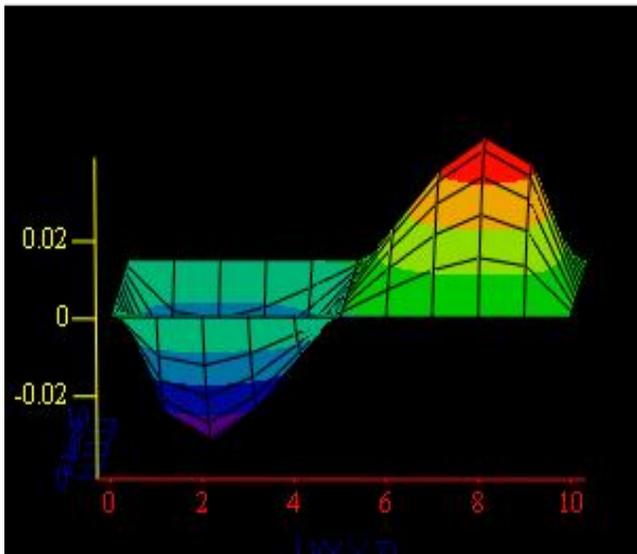


Figure 3: The graph of the distribution of the function $u(x,y,z,t)$ in the plane XOY at $z=0.3$ and $t=0.09$

Table 4: Displacement values $v(x,y,z,t)$ at $z=0.3$ and $t=0.09$

| x | y | 0 | 0.1 | 0.2 | 0.3 | 0.4 | 0.5 |
|-----|-----|----------|----------|----------|----------|---------|---------|
| 0 | 0 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 |
| 0.1 | 0 | -0.01135 | -0.01903 | -0.02568 | -0.03013 | 0.00000 | 0.00000 |
| 0.2 | 0 | -0.01362 | -0.02337 | -0.03165 | -0.03713 | 0.00000 | 0.00000 |
| 0.3 | 0 | -0.01066 | -0.01834 | -0.02486 | -0.02918 | 0.00000 | 0.00000 |
| 0.4 | 0 | -0.00567 | -0.00977 | -0.01324 | -0.01554 | 0.00000 | 0.00000 |
| 0.5 | 0 | -0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 |
| 0.6 | 0 | 0.00567 | 0.00977 | 0.01324 | 0.01554 | 0.00000 | 0.00000 |
| 0.7 | 0 | 0.01066 | 0.01834 | 0.02486 | 0.02918 | 0.00000 | 0.00000 |
| 0.8 | 0 | 0.01362 | 0.02337 | 0.03165 | 0.03713 | 0.00000 | 0.00000 |
| 0.9 | 0 | 0.01135 | 0.01903 | 0.02568 | 0.03013 | 0.00000 | 0.00000 |
| 1 | 0 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 |

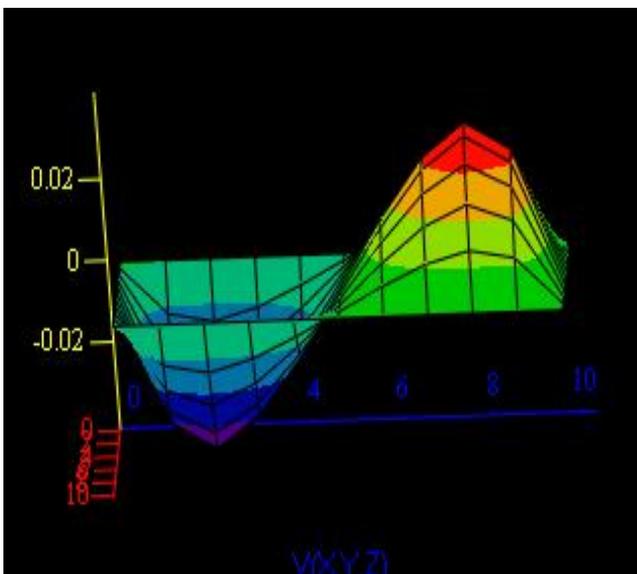


Figure 4: The graph of the distribution of the function $v(x,y,z,t)$ in the plane XOY at $z=0.3$ and $t=0.09$

Table 5: Displacement values $w(x,y,z,t)$ at $x=0.3$ and $t=0.09$

| x | z | 0 | 0.1 | 0.2 | 0.3 | 0.4 | 0.5 |
|-----|-----|----------|----------|----------|----------|----------|----------|
| 0 | 0 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 |
| 0.1 | 0 | -0.01695 | -0.02967 | -0.04034 | -0.04735 | -0.04978 | -0.04978 |
| 0.2 | 0 | -0.01983 | -0.03530 | -0.04812 | -0.05651 | -0.05941 | -0.05941 |
| 0.3 | 0 | -0.01558 | -0.02775 | -0.03785 | -0.04445 | -0.04673 | -0.04673 |
| 0.4 | 0 | -0.00831 | -0.01481 | -0.02019 | -0.02372 | -0.02493 | -0.02493 |
| 0.5 | 0 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 |
| 0.6 | 0 | 0.00831 | 0.01481 | 0.02019 | 0.02372 | 0.02493 | 0.02493 |
| 0.7 | 0 | 0.01558 | 0.02775 | 0.03785 | 0.04445 | 0.04673 | 0.04673 |
| 0.8 | 0 | 0.01983 | 0.03530 | 0.04812 | 0.05651 | 0.05941 | 0.05941 |
| 0.9 | 0 | 0.01695 | 0.02967 | 0.04034 | 0.04735 | 0.04978 | 0.04978 |
| 1 | 0 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 |

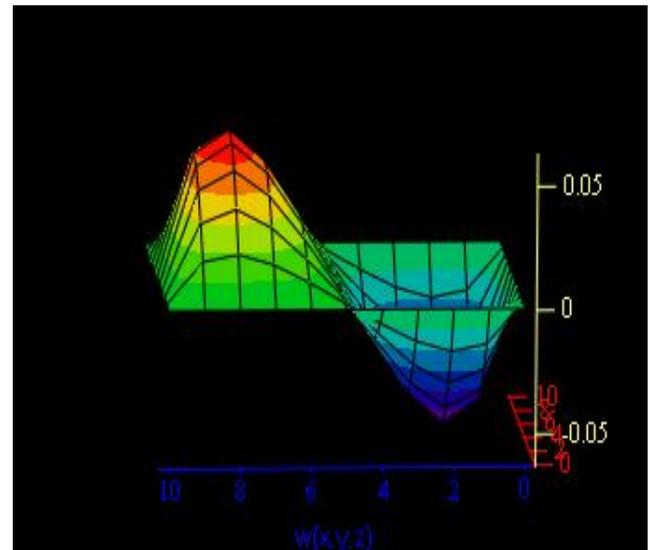


Figure 5: The graph of the distribution of the function $w(x,y,z,t)$ in the plane ZOY at $x=0.3$ and $t=0.09$

Table 6: Temperature values $T(x,y,z,t)$ at $z=0.5$ and $t=0.05$

| x | z | 0 | 0.1 | 0.2 | 0.3 | 0.4 | 0.5 |
|-----|-----|---------|----------|----------|----------|----------|----------|
| 0 | 0 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 |
| 0.1 | 0 | 1.98896 | 3.63207 | 4.98561 | 5.86017 | 6.16171 | 6.16171 |
| 0.2 | 0 | 3.63200 | 6.63959 | 9.11426 | 10.71307 | 11.26434 | 11.26434 |
| 0.3 | 0 | 4.98561 | 9.11433 | 12.51135 | 14.70607 | 15.46281 | 15.46281 |
| 0.4 | 0 | 5.86019 | 10.71313 | 14.70607 | 17.28578 | 18.17527 | 18.17527 |
| 0.5 | 0 | 6.16175 | 11.26439 | 15.46281 | 18.17527 | 19.11052 | 19.11052 |
| 0.6 | 0 | 5.86019 | 10.71310 | 14.70607 | 17.28578 | 18.17527 | 18.17527 |
| 0.7 | 0 | 4.98561 | 9.11427 | 12.51135 | 14.70607 | 15.46281 | 15.46281 |
| 0.8 | 0 | 3.63200 | 6.63949 | 9.11426 | 10.71307 | 11.26434 | 11.26434 |
| 0.9 | 0 | 1.98893 | 3.63200 | 4.98556 | 5.86015 | 6.16171 | 6.16171 |
| 1 | 0 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 | 0.00000 |

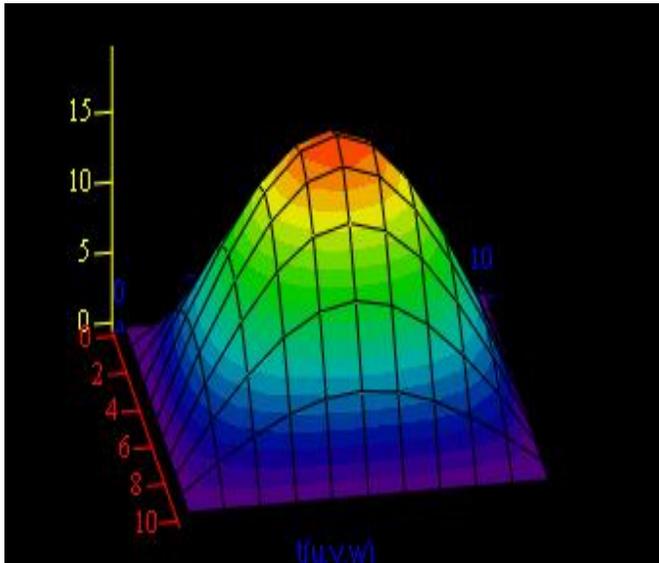


Figure 6: Temperature distribution graph of $T(x, y, z, t)$ in the plane XOY at $z=0.5$ and $t=0.05$

Note that, in the considered model problem, the axis of rotation was the OZ axis. Tables (1-2) and Figures (1-2) show the appearance of plastic zones, and Table 6 and Figure 6 show the temperature field. Since the initial and boundary conditions are set symmetrically, in Tables (3-5) and Figures (3-5) you can see the symmetry of the displacements $u(x,y,z)$ and $v(x,y,z)$, and displacements $w(x,y,z)$ are different from others. This shows the effect of anisotropy on the elastic-plastic state of a transversely isotropic parallelepiped and the reliability of the numerical results obtained.

5. CONCLUSION

A coupled dynamic thermoplastic boundary value problem for the deformation theory of transversely isotropic bodies is formulated. A discrete analogue of the problem is compiled by the finite difference method. On the basis of explicit finite-difference equations, recurrent formulas are obtained that allow calculating the numerical values of the required functions. A numerically dynamic coupled first boundary value problem on a thermoplastic transversely isotropic parallelepiped is solved. 3D-graphs of the distribution of displacement functions $u(x,y,z)$, $v(x,y,z)$, $w(x,y,z)$ and temperatures $T(x,y,z)$. The propagation of plasticity zones in various sections of a transversely isotropic parallelepiped under the action of a temperature field is investigated.

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